# Deeply Virtual Compton Scattering @ JLab

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# On behalf of the Hall A and CLAS collaborations

From GPDs to DVCS, to GPDs back Onto the DVCS harmonic structure E00-110 experiment in Hall A Scaling tests & GPD measurement E1-DVCS experiment at CLAS in Hall B Summary



# GPDs from Theory to Experiment



2. The GPDs enter the DVCS amplitude as an integral over x:

- GPDs appear in the real part through a PP integral over x
- GPDs appear in the imaginary part but at the line x= $\xi$

$$T^{DVCS} = \int_{-1}^{+1} \frac{GPD(x,\xi,t)}{x-\xi+i\varepsilon} dx + \cdots$$
$$= P \int_{-1}^{+1} \frac{GPD(x,\xi,t)}{x-\xi} dx - i\pi GPD(x=\xi,\xi,t) + \cdots$$

#### Experimental observables linked to GPDs

#### 3. Experimentally, DVCS is undistinguishable with Bethe-Heitler



However, we know FF at low t and BH is fully calculable

Using a polarized beam on an unpolarized target, 2 observables can be measured:

$$\frac{d^4\sigma}{dx_B dQ^2 dt d\varphi} \approx \left|T^{BH}\right|^2 + 2T^{BH} \cdot \operatorname{Re}\left(T^{DVCS}\right) + \left|T^{DVCS}\right|^2$$

$$\frac{d^{4} \overrightarrow{\sigma} - d^{4} \overleftarrow{\sigma}}{dx_{B} dQ^{2} dt d\varphi} \approx 2T^{BH} \cdot \operatorname{Im}(T^{DVCS}) + \left[ \left| T^{DVCS} \right|^{2} - \left| T^{DVCS} \right|^{2} \right]$$

$$A^{\dagger} JLab \text{ energies,}$$

$$|T^{DVCS}|^{2} \text{ should be small}$$

Kroll, Guichon, Diehl, Pire, ...



Into the harmonic structure of DVCS



Belitsky, Mueller, Kirchner

$$\frac{d^{4}\sigma}{dx_{B}dQ^{2}dtd\varphi} = \frac{1}{\Pr_{1}(\varphi)\Pr_{2}(\varphi)}\Gamma_{1}(x_{B},Q^{2},t)\left\{c_{0}^{BH}+c_{1}^{BH}\cos\varphi+c_{2}^{BH}\cos2\varphi\right\}$$
$$+\frac{1}{\Pr_{1}(\varphi)\Pr_{2}(\varphi)}\Gamma_{2}(x_{B},Q^{2},t)\left\{\frac{c_{0}^{I}+c_{1}^{I}\cos\varphi}{c_{0}^{I}+c_{2}^{I}\cos2\varphi+c_{3}^{I}\cos3\varphi}\right\}$$

$$\frac{d^4 \overrightarrow{\sigma} - d^4 \overleftarrow{\sigma}}{dx_B dQ^2 dt d\varphi} = \frac{\Gamma(x_B, Q^2, t)}{\mathbf{P}_1(\varphi) \mathbf{P}_2(\varphi)} \left\{ \frac{s_1^I \sin \varphi}{s_1^I \sin \varphi} + s_2^I \sin 2\varphi \right\}$$

1. Twist-2 terms should dominate  $\sigma$  and  $\Delta\sigma$ 

2. All coefficients have  $Q^2$  dependence which can be tested!

<u>Re-stating the problem (difference of cross-section)</u>:

$$\frac{d^{4} \stackrel{\rightarrow}{\sigma} - d^{4} \stackrel{\leftarrow}{\sigma}}{dx_{B} dQ^{2} dt d\varphi} = \frac{\Gamma(x_{B}, Q^{2}, t)}{P_{1}(\varphi)P_{2}(\varphi)} \left\{ s_{1}^{\prime} \sin \varphi + s_{2}^{\prime} \sin 2\varphi \right\}$$

$$s_{1}^{\prime} = 8Ky(2 - y) \left[ \text{Im } C^{\prime}(F) \right] \quad \longleftarrow \quad \text{What we measure}$$

$$C^{\prime}(F) = F_{1} H + \frac{x_{B}}{2 - x_{B}} (F_{1} + F_{2}) \tilde{H} - \frac{t}{4M^{2}} F_{2} E$$

$$\text{Im } H = \pi \sum_{q} e_{q}^{2} \left\{ H^{q}(\xi, \xi, t) - H^{q}(-\xi, \xi, t) \right\}$$

$$\text{GPD III}$$

The asymmetry can be written as:

$$\frac{d^{4} \overrightarrow{\sigma} - d^{4} \overleftarrow{\sigma}}{d^{4} \overrightarrow{\sigma} + d^{4} \overleftarrow{\sigma}} = \Gamma_{A} (x_{B}, Q^{2}, t) \frac{s_{1}^{I} \sin \varphi + s_{2}^{I} \sin 2\varphi}{c_{0}^{I} + c_{0}^{BH} + (c_{1}^{I} + c_{1}^{BH}) \cos \varphi + \dots}$$

Pros: easier experimentally, smaller RC

<u>Cons:</u> - extraction of GPDs model-dependent (denominator complicated and not well known) - Large effects of the BH propagators in the denominator

Asymmetries are largely used in CLAS and HERMES measurements, where acceptance and systematics are more difficult to estimate.

# E00-110 experimental setup and performances



# E00-110 kinematics

Kin	$Q^2$	$x_B$	$\theta_{\gamma^*}$	W
	$({\sf GeV}^2)$		(deg.)	(GeV)
1	1.5	0.36	22.3	1.9
2	1.9	0.36	18.3	2.0
3	2.3	0.36	14.8	2.2

The calorimeter is centered on the virtual photon direction



50 days of beam time in the fall 2004, at 2.5µA intensity  $\int Lu \cdot dt = 13294 \text{ fb}^{-1}$ 

HRS: Cerenkov, vertex, flat-acceptance cut with R-functions

<u>Calo:</u> 1 cluster in coincidence in the calorimeter above 1 GeV



<u>With both</u>: subtract accidentals, build missing mass of  $(e,\gamma)$  system

Using  $\pi^0 \rightarrow 2\gamma$  events in the calorimeter, the  $\pi^0$  contribution is subtracted bin by bin



#### Analysis - Exclusivity check using Proton Array and MC

Using Proton-Array, we compare the missing mass spectrum of the triple and double-coincidence events.

The missing mass spectrum using the Monte-Carlo gives the same position and width. Using the cut shown on the Fig., the contamination from inelastic channels is estimated to be under 3%.





#### Difference of cross-sections





Twist 4+ contributions are smaller than 10%

## Total cross-section



# E1-DVCS @ CLAS : a dedicated DVCS experiment



#### E1-DVCS kinematical coverage and binning





Remaining  $\pi^0$  contamination up to 20%, subtracted bin by bin using p0 events and MC estimation of  $\pi^0(1\gamma)$  to  $\pi^0(2\gamma)$  acceptance ratio

CLAS results were very preliminary and cannot be put on the web.

# Summary

Cross-section difference (Hall A):

□High statistics test of scaling: Strong support for twist-2 dominance

□ First model-independent extraction of GPD linear combination from DVCS data in the twist-3 approximation

□ Upper limit set on twist-4+ effects in the cross-section difference: twist>3 contribution is smaller than 10%

Total cross-section (Hall A):

□ Bethe-Heitler is not dominant everywhere

 $\square$  |DVCS|<sup>2</sup> terms might be sizeable but almost impossible to extract using only total cross-section: e<sup>+</sup>/e<sup>-</sup> or  $\mu^+/\mu^-$  beams seem necessary

Despite this, we performed a measurement of 2 different GPD integrals

BSA (CLAS):

□ Preliminary data in large kinematic range and good statistics !