

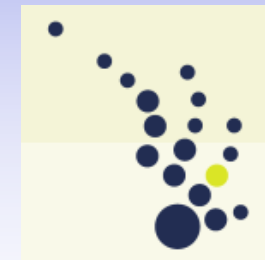
A Finite Element Model for the Thermal Transport in Solid Targets



A. Raccanelli

R. Krause

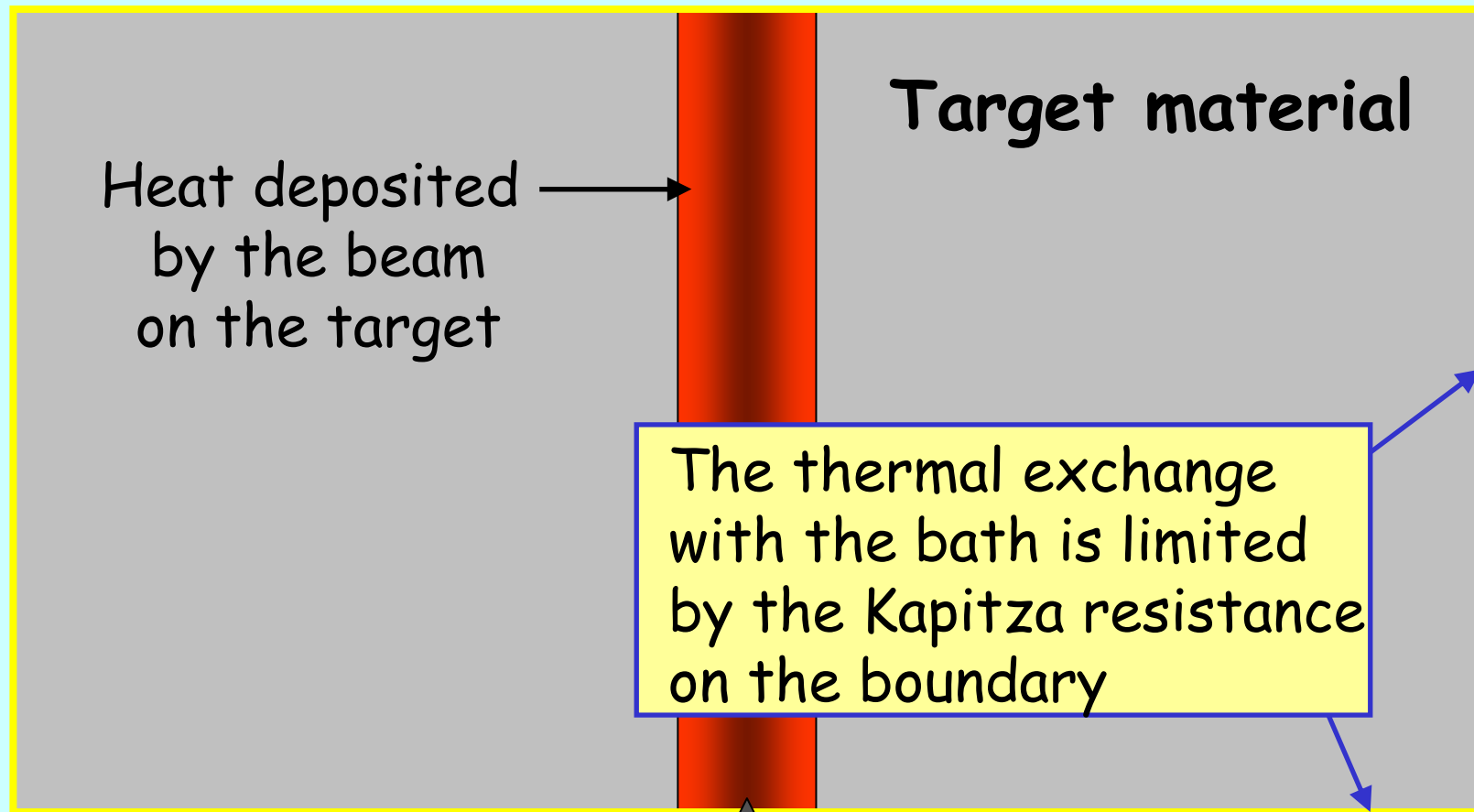
H. Dutz



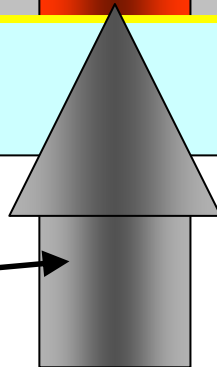
INS-BONN



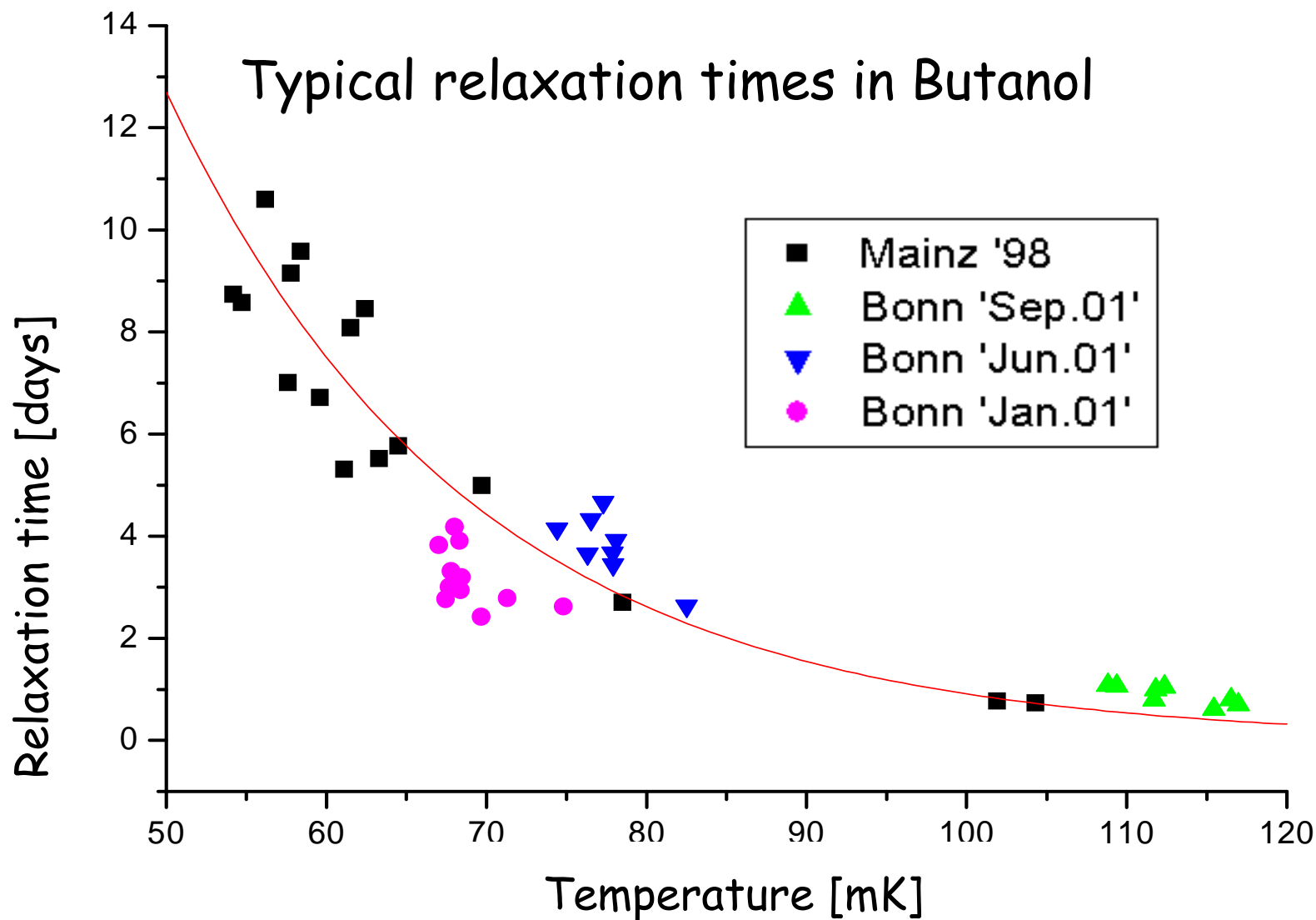
^3He - ^4He mixture, $T=T_0$

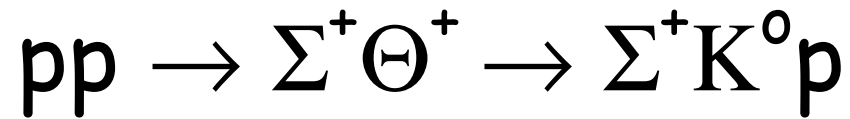


Proton beam



Temperature affects polarization





Requirements:

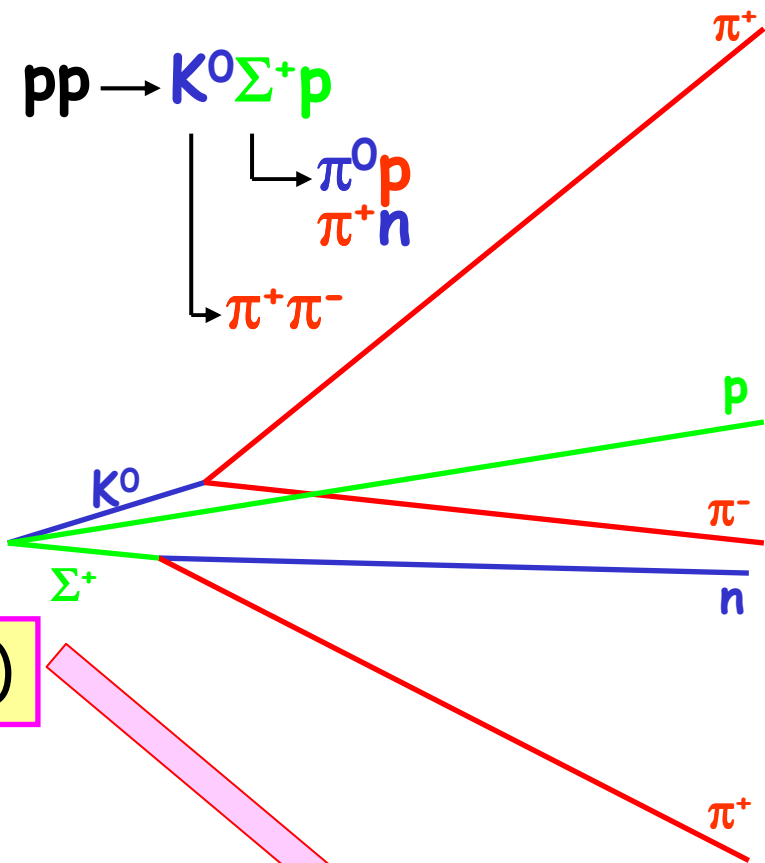
★ For detection statistics:

high beam intensity ($N=10^7/s$)

★ For accuracy of vertex reconstruction:

high beam collimation ($\sim 1\text{mm}^2$)

detection of Σ^+ ($c\tau \approx 2.4\text{cm}$)



Target design

Frozen spin mode

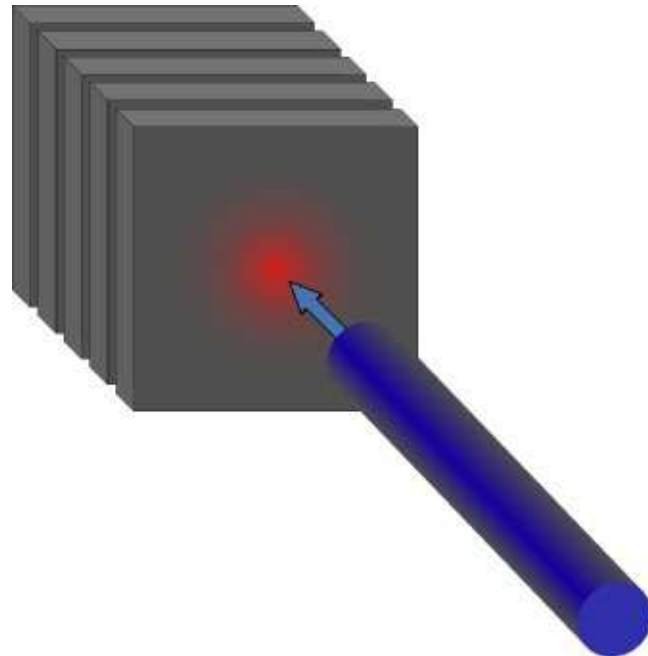
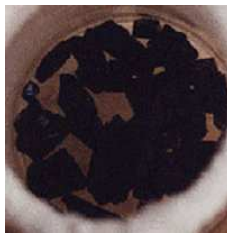
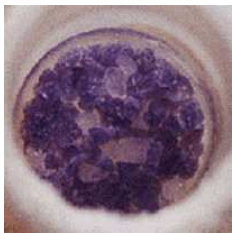
The Kapitza resistance

From the acoustic mismatch theory

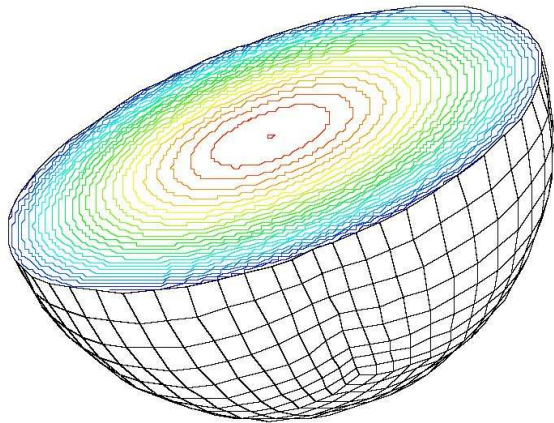
$$\frac{\dot{Q}}{A} = \frac{\pi^2 k_B^4 T^4 \rho_h v_h}{30 \hbar \rho_s v_s^3}$$

$$\frac{\dot{Q}}{A} = \alpha (T^4 - T_0^4)$$

Heat exchanged between the solid and the bath at $T=T_0$



An analytical solution exists only for a sphere with thermal conductivity $k = \text{const.}$



$$-k\nabla^2 T = \dot{Q}_V \quad (k = \text{const.})$$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) T(r) + \frac{\dot{Q}_V}{k} = 0$$

$$T(r) = -\frac{\dot{Q}_V}{6k} \cdot r^2 + \text{const.}$$

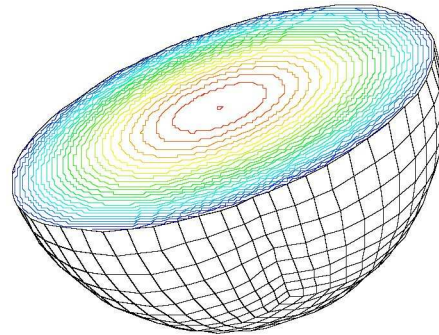
$$\Delta T = T(0) - T(R) = \frac{\dot{Q}_V}{6k} \cdot R^2$$

On the surface:

$$\frac{\dot{Q}}{A} = \alpha(T^4 - T_0^4)$$

The sphere is the „worste case“

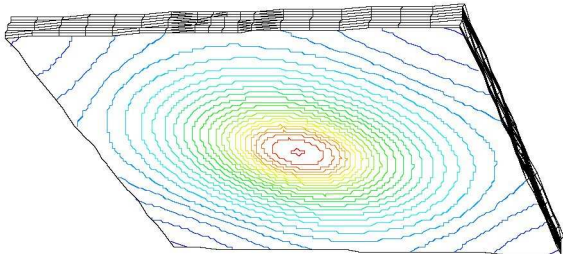
$$\frac{\dot{Q}}{A} = \alpha (T^4 - T_0^4)$$



$$T = \sqrt[4]{\frac{\dot{Q}}{A \cdot \alpha} + T_0^4}$$

$$T = \sqrt[4]{\frac{\dot{Q}_v \cdot \frac{4}{3} \pi R^3}{\alpha \cdot 4 \pi R^2} + T_0^4}$$

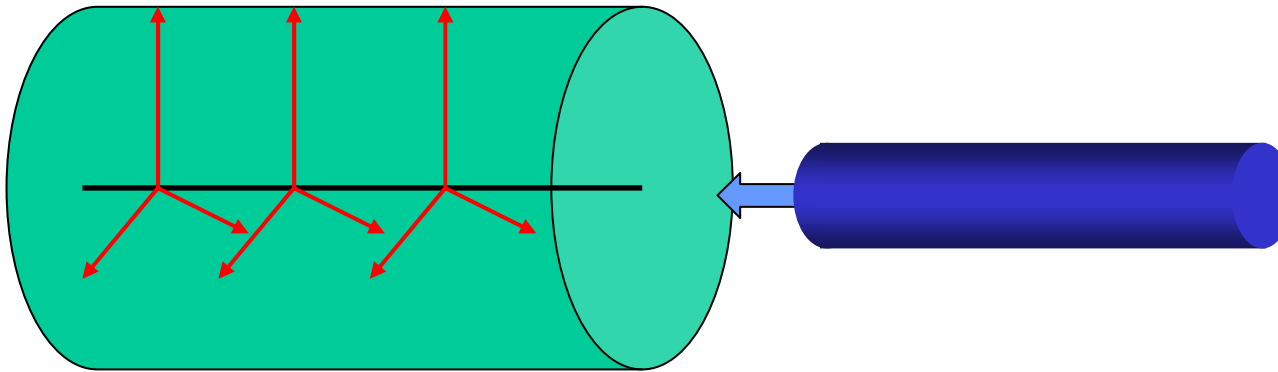
$$T = \sqrt[4]{\frac{\dot{Q}_v}{3\alpha} R + T_0^4}$$



Thickness decreases, area \approx constant

Other common assumptions:

- 1) The heat is dissipated through the lateral surface only

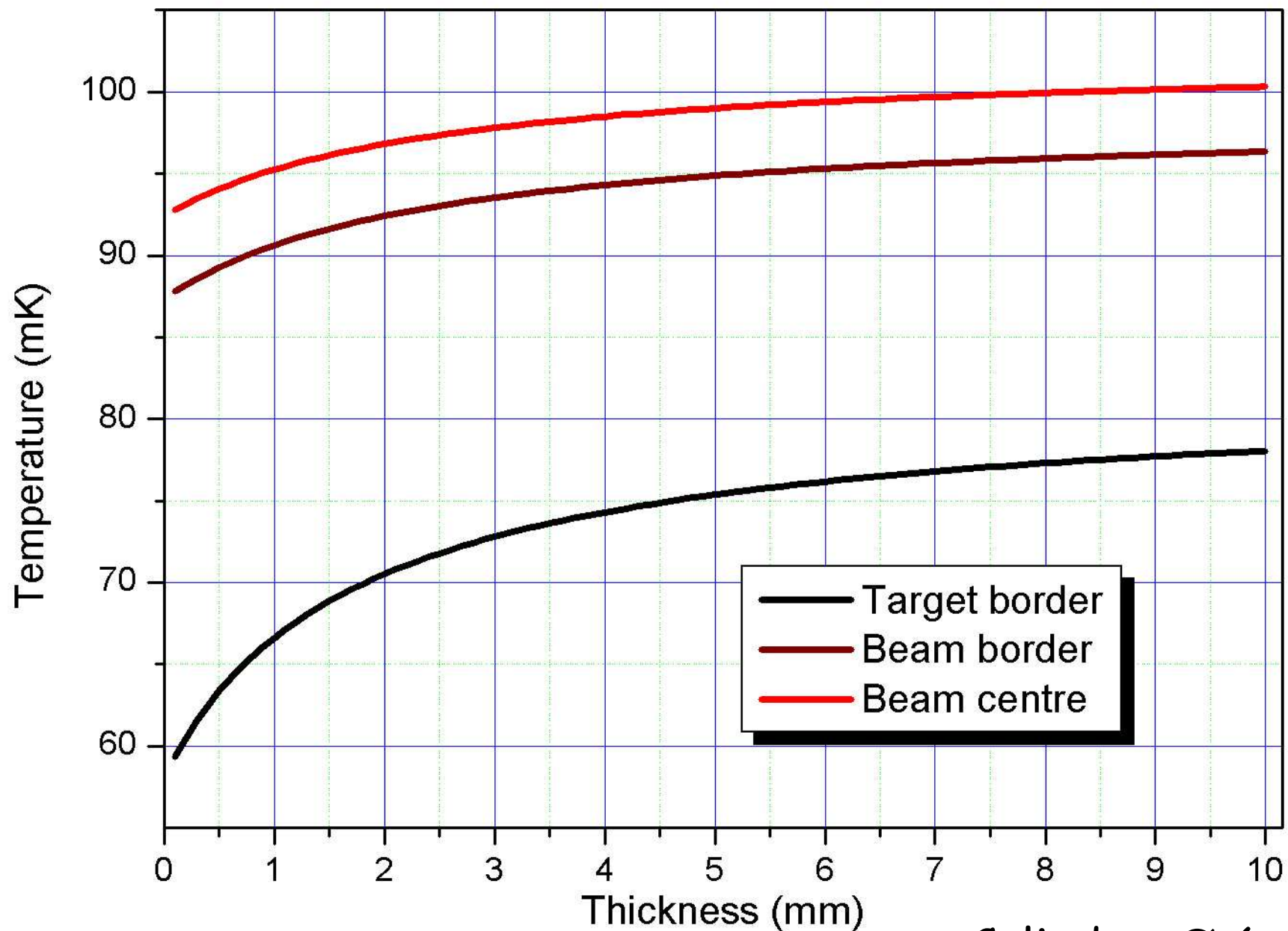


- 2) The Kapitza "jump" is constant on the surface:

$$\frac{\dot{Q}}{A} = \alpha(T^4 - T_0^4)$$

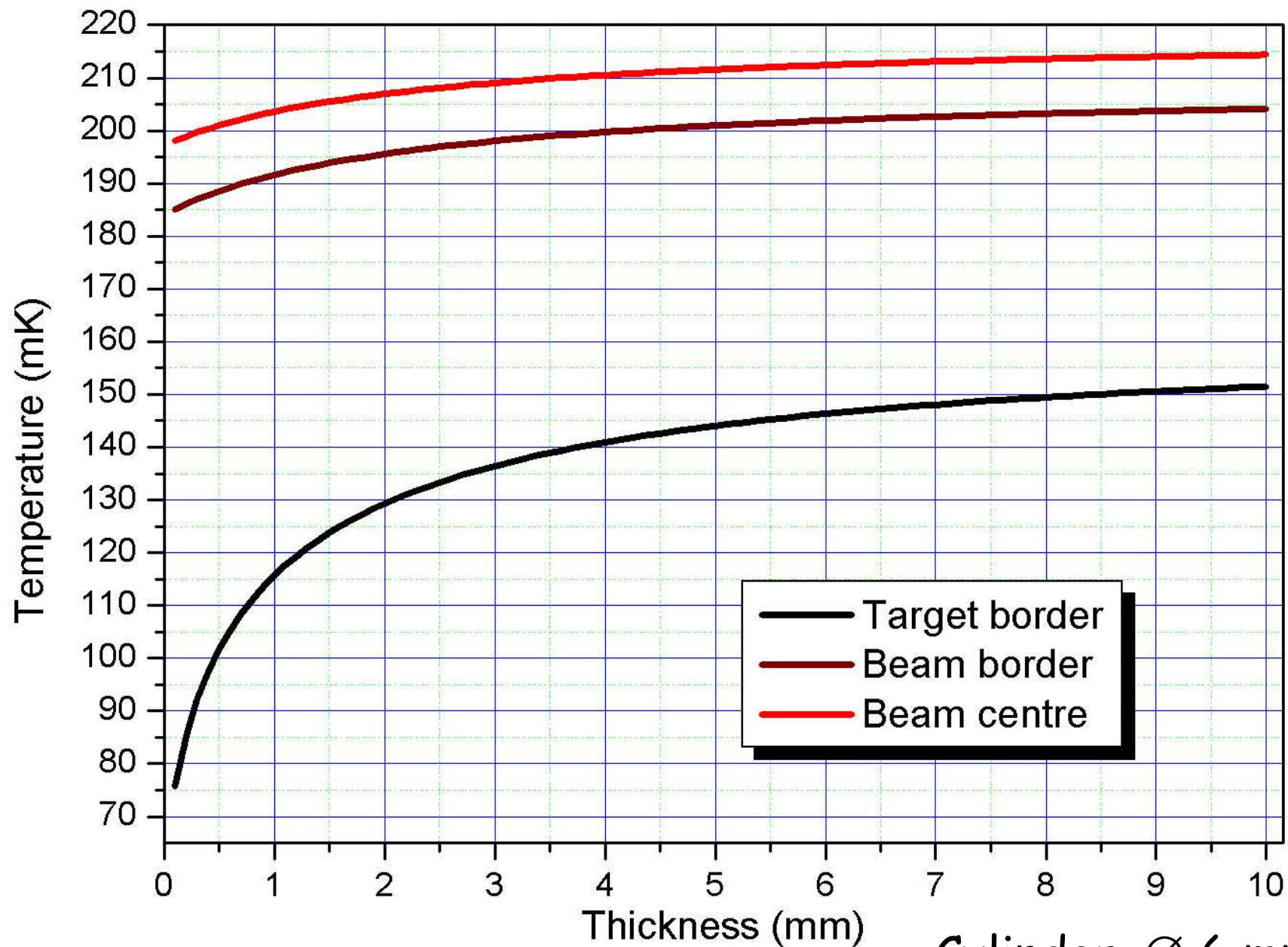
$$T = \sqrt[4]{\frac{\dot{Q}}{A \cdot \alpha} + T_0^4}$$

$N = 500 \text{ kHz}$, beam radius = 0.56 mm



Cylinder, $\varnothing 6 \text{ mm}$

$N = 10\text{MHz}$, beam radius = 0.56 mm



Cylinder, \varnothing 6 mm

The model - (1) Heat-conduction equation

The time-independent heat-conduction equation

$$-k\nabla^2 T = -k \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) = q \quad (\text{Poisson eq.})$$

for a non constant thermal conductivity, k , is:

$$-\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left(k \frac{\partial T}{\partial z} \right) = q$$

$$k = k(T) = aT^\beta \quad \text{with } a \text{ and } \beta \text{ free parameters}$$

The model - (2) The Kapitza resistance

$$\dot{Q} = -kA \frac{\partial T}{\partial n}$$

+

$$\text{Robin boundary condition} \\ \frac{\partial T}{\partial n} + h \cdot T = f$$

$$\frac{\dot{Q}}{A} = \alpha(T^4 - T_0^4)$$



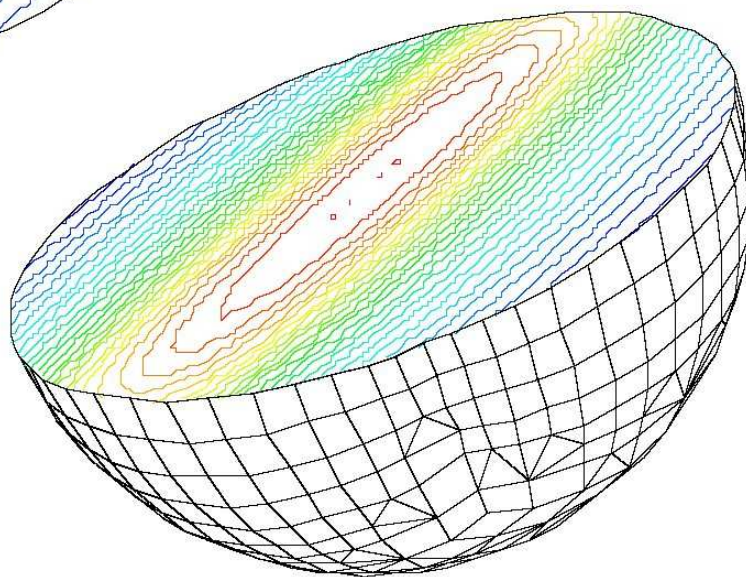
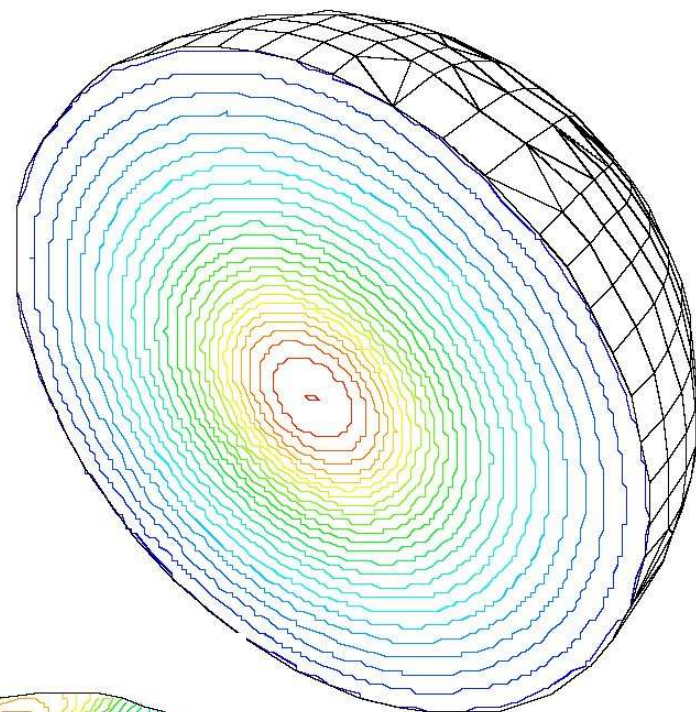
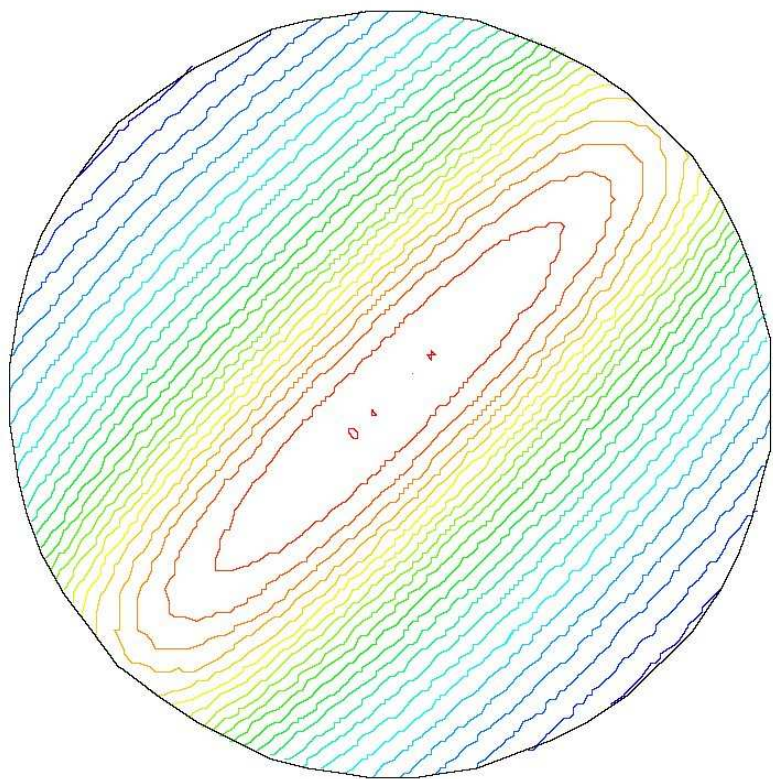
$$k \frac{\partial T}{\partial n} + \alpha \cdot T^4 = \alpha \cdot T_0^4$$

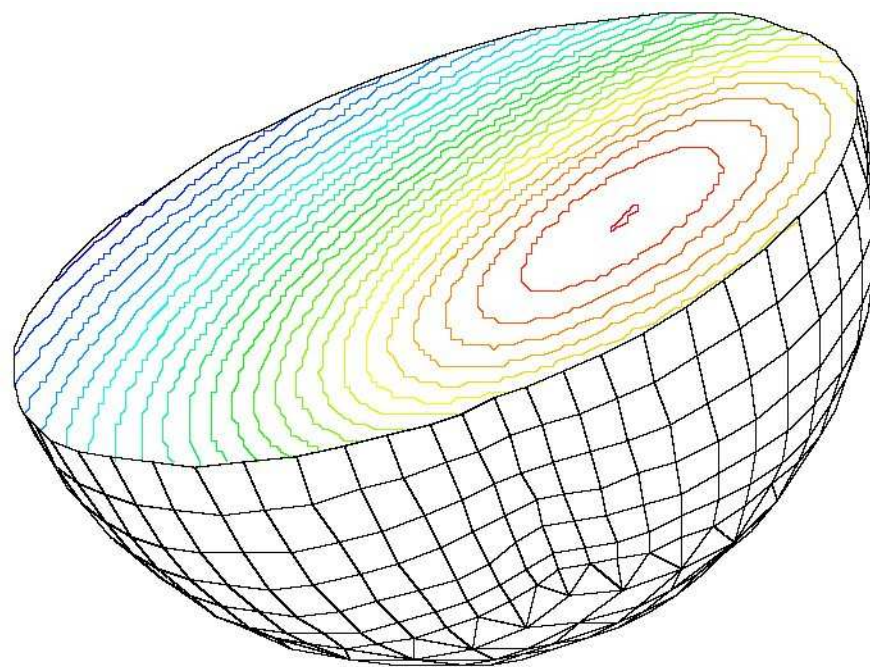
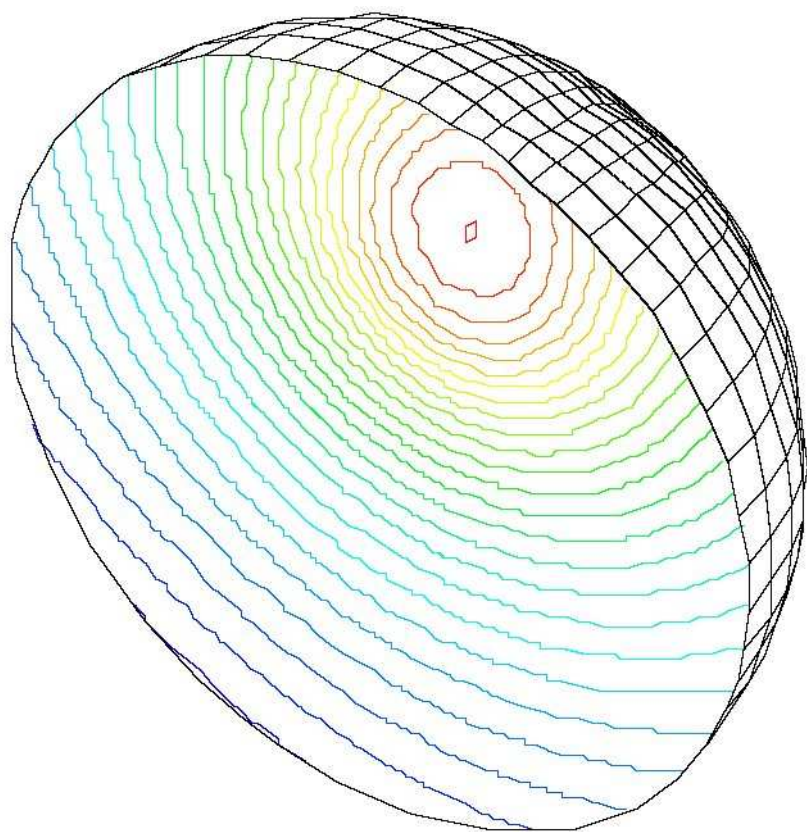
The implementation of the model is based on the UG package.

P. Bastian et al., Comput. visual. Sci. 1(1997)

The solution of the linear system is based on the multigrid method.

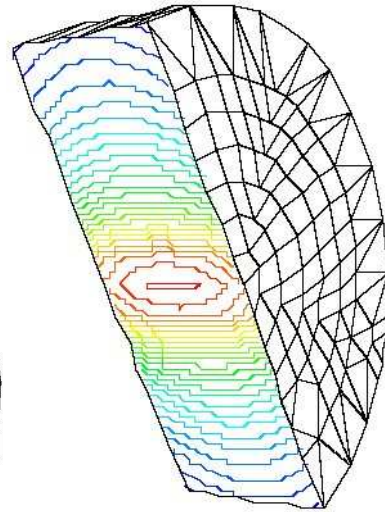
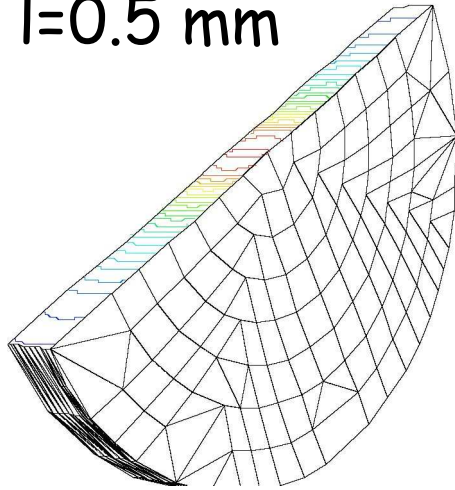
+ adaptive grids



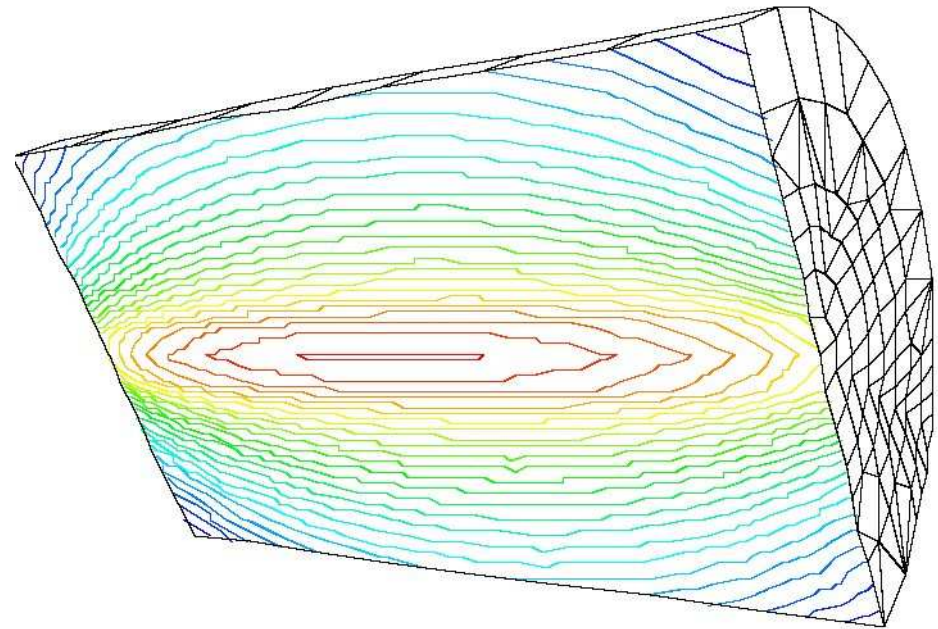
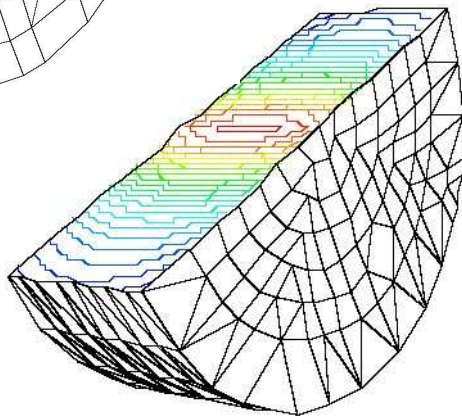


Cylinder, \varnothing 6 mm

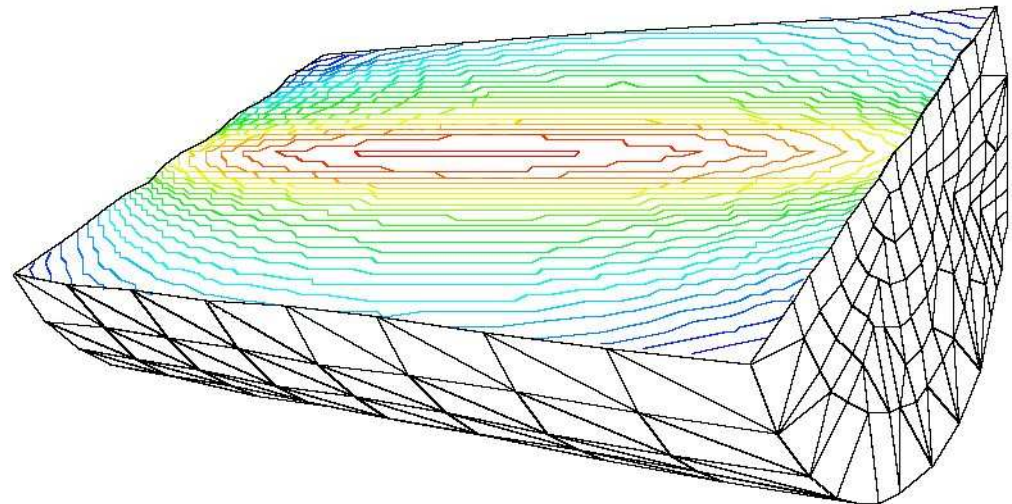
$l=0.5$ mm

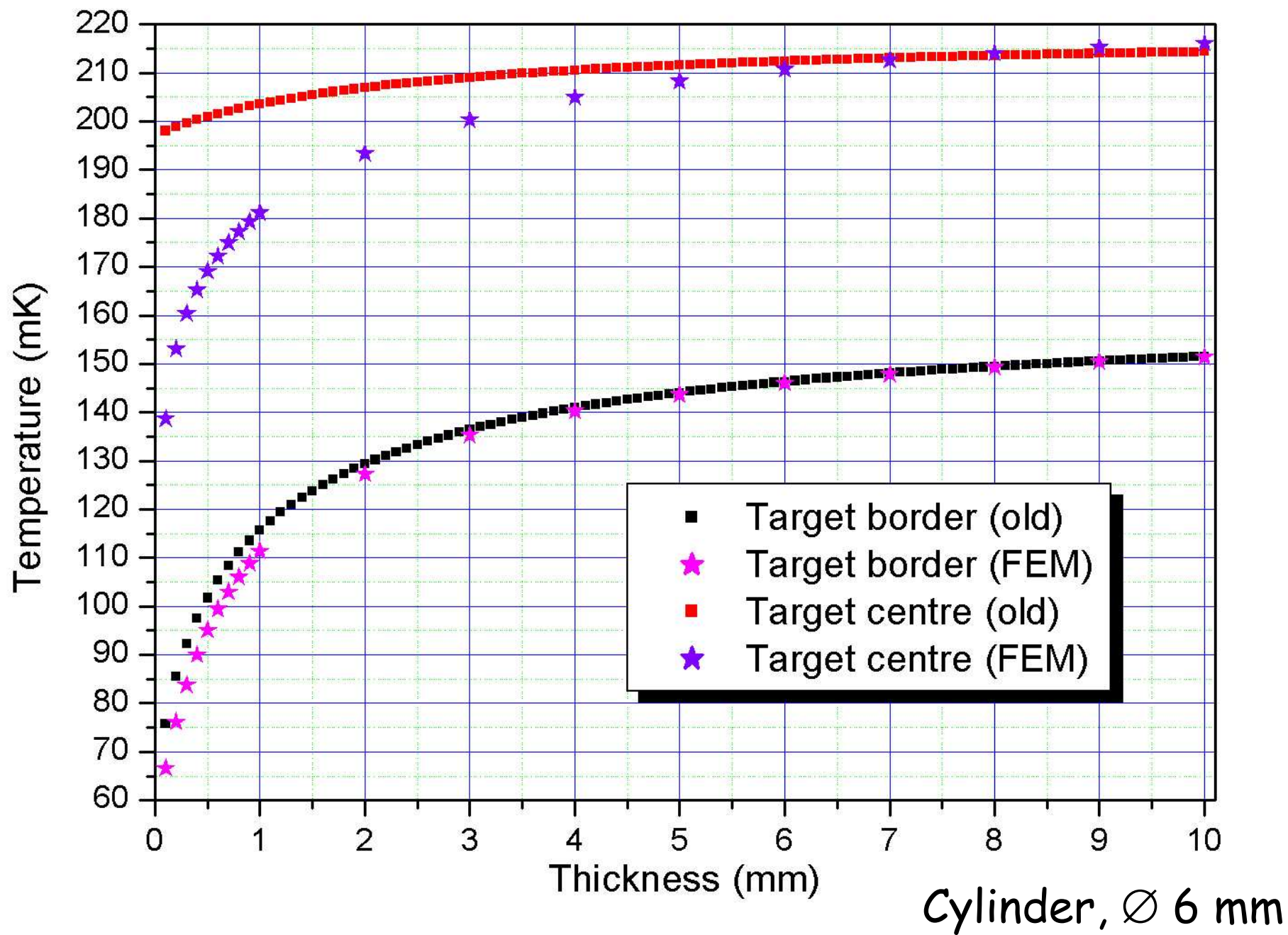


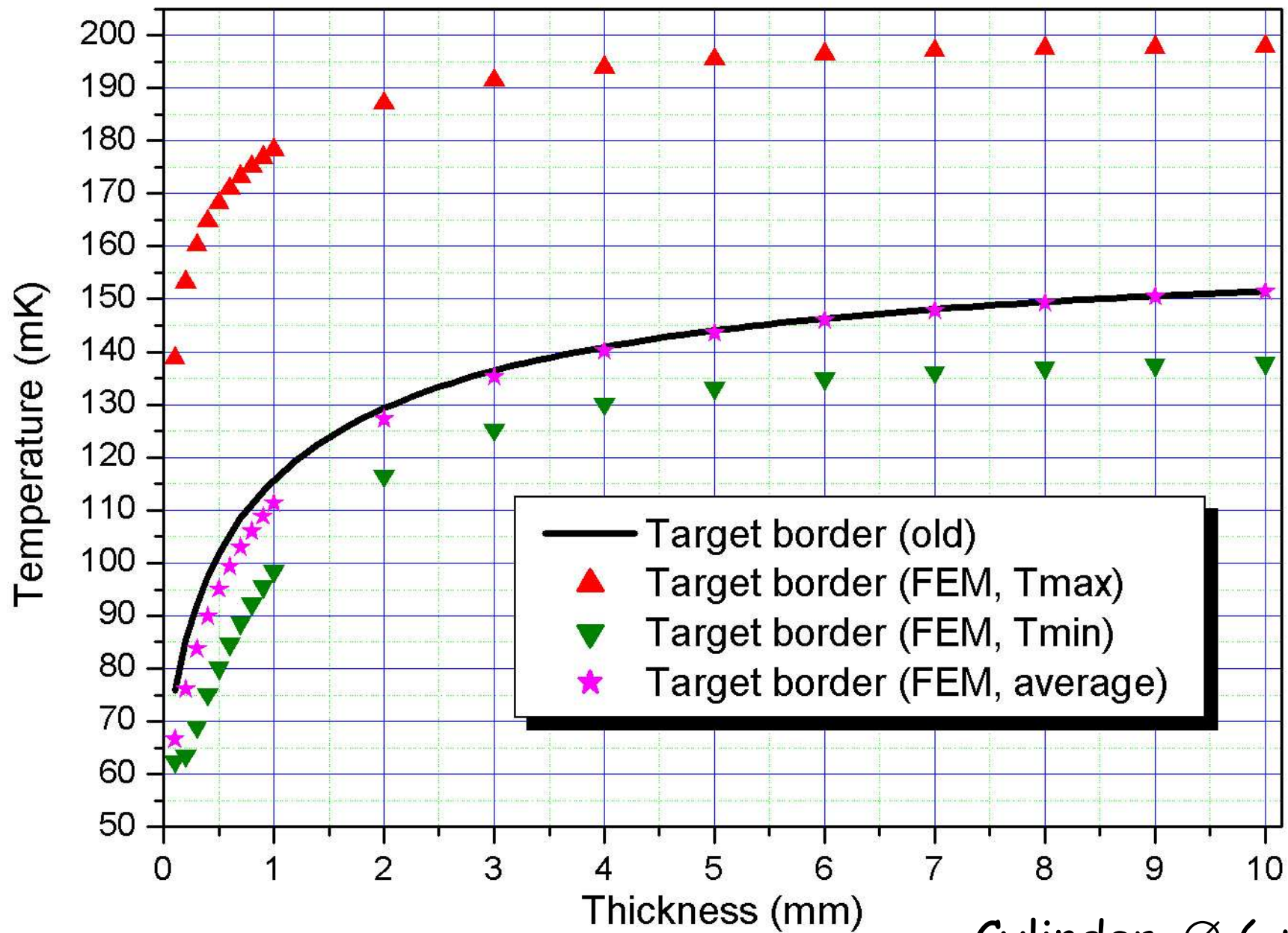
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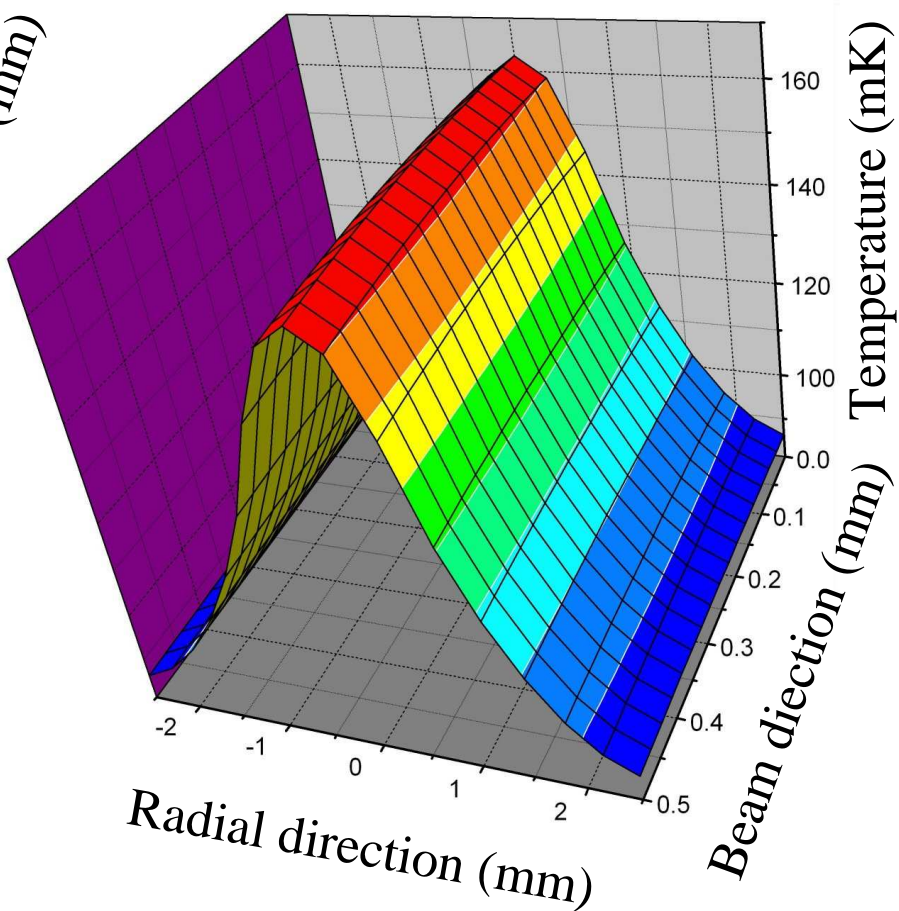
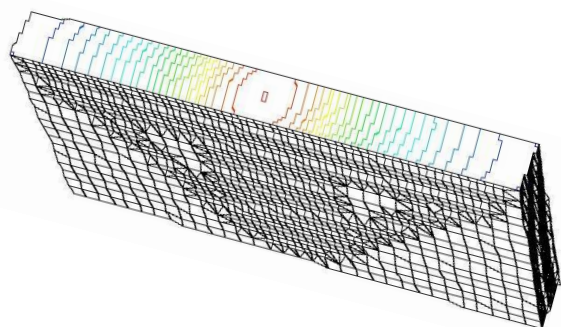
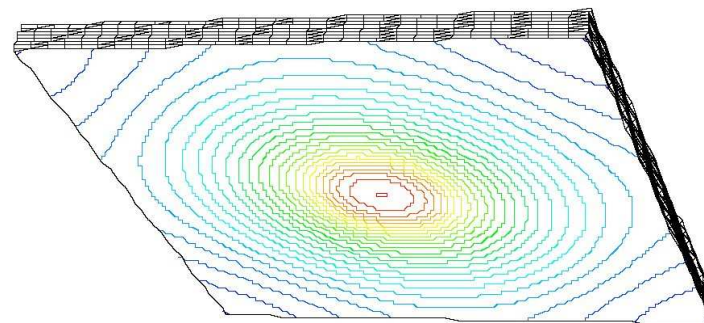
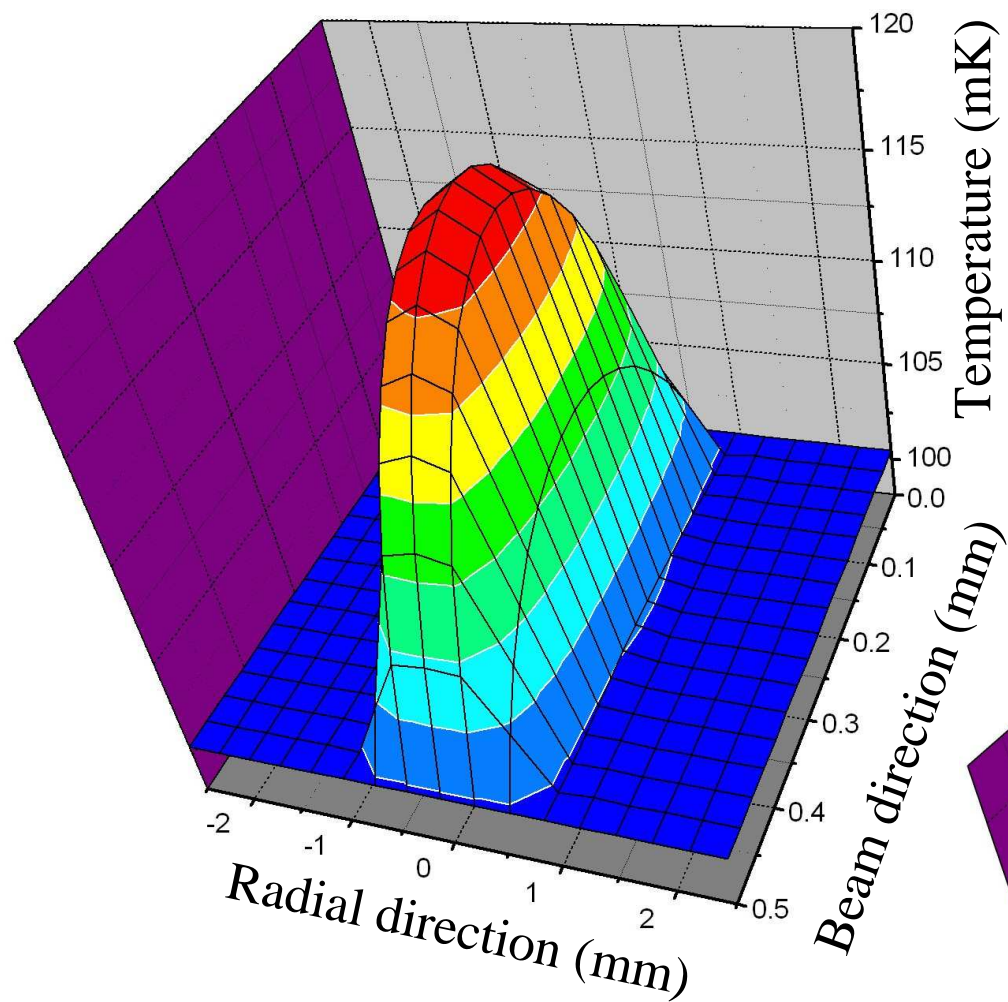


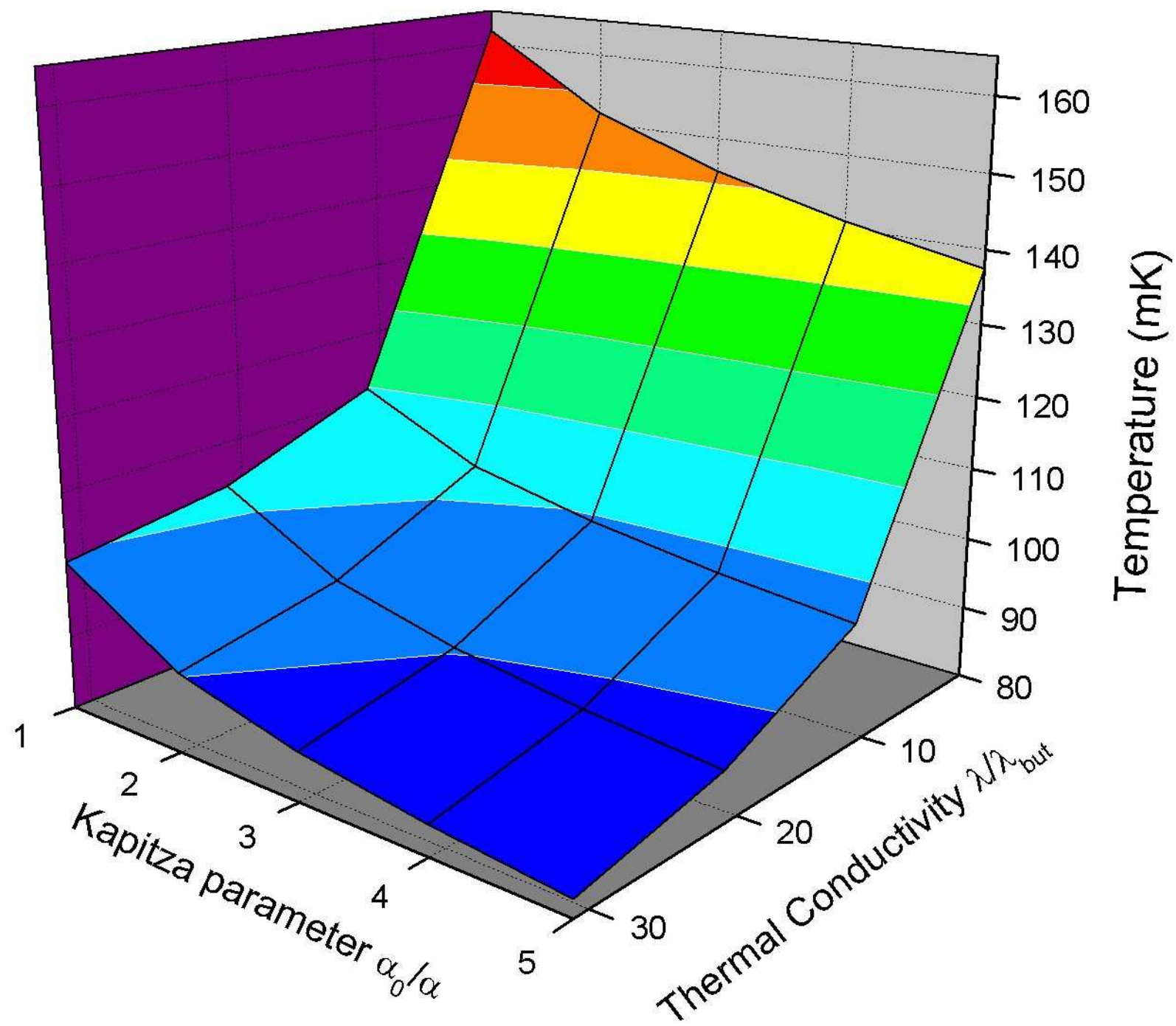
$l=10$ mm

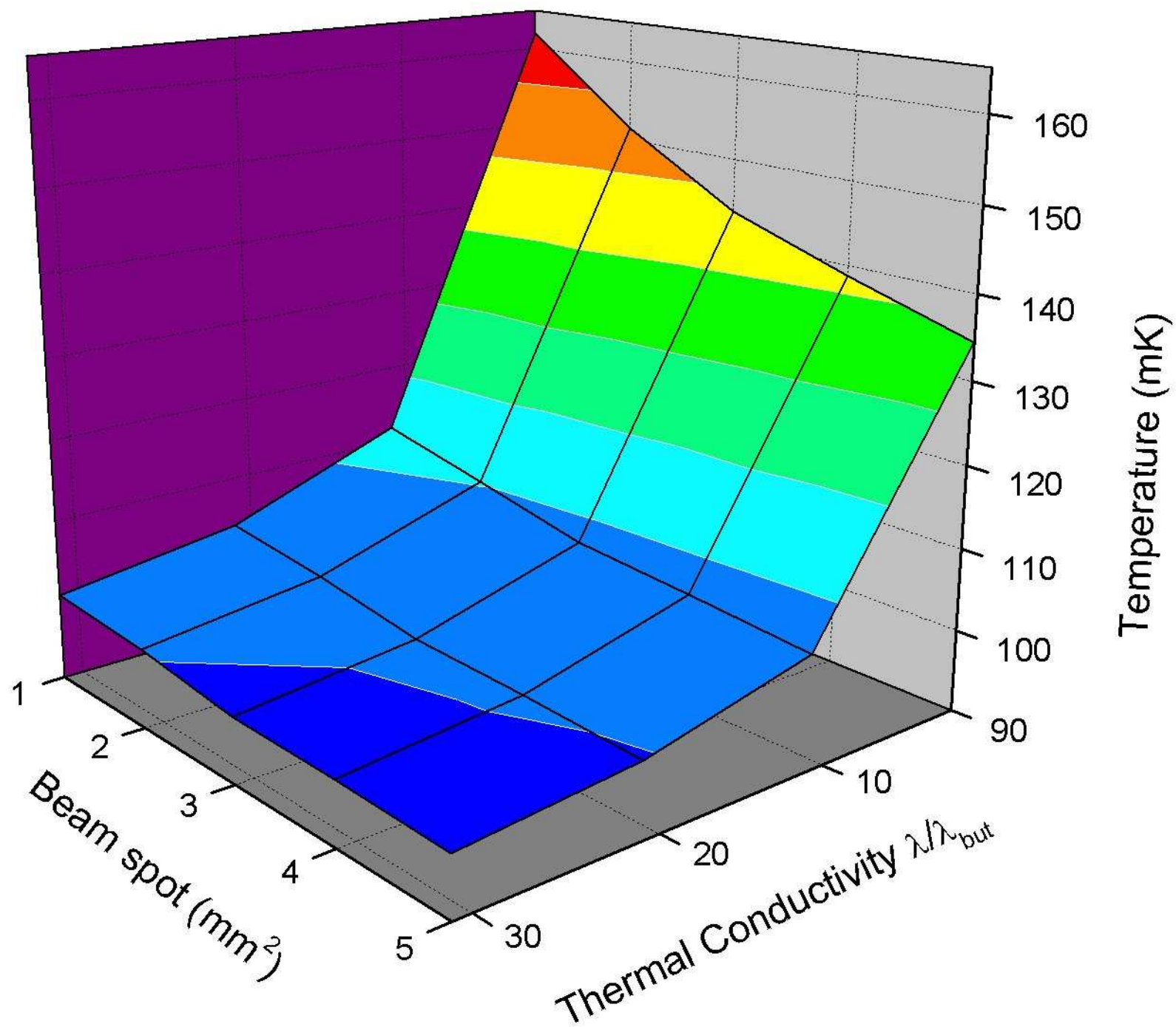


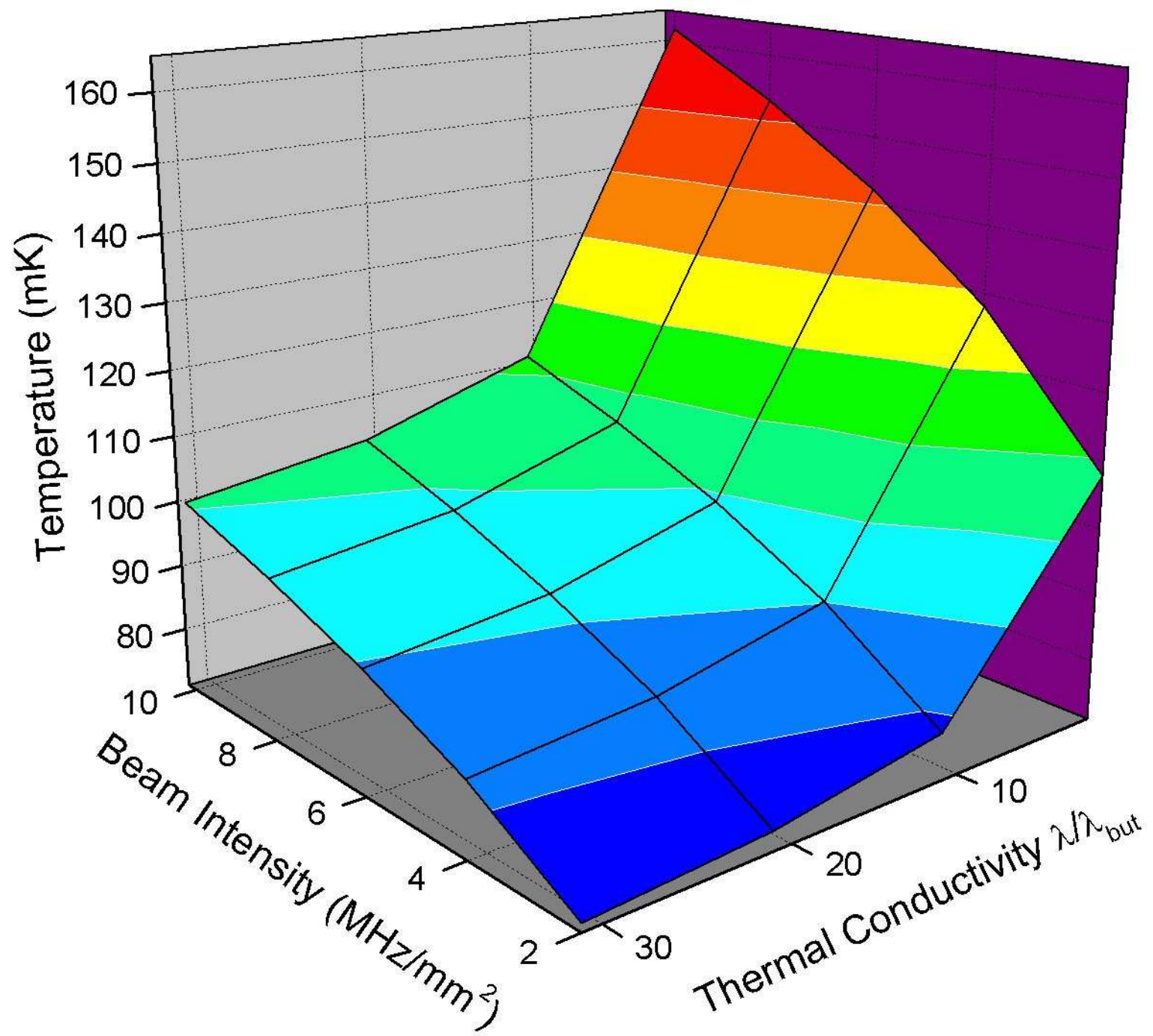












CONCLUSIONS

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THANK YOU!